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# James B. Miller<sup>\*</sup>, Andrew J. Gellman

National Energy Technology Laboratory, United States Department of Energy, Pittsburgh, PA 15236, USA Department of Chemical Engineering, Carnegie Mellon University, Pittsburgh, PA 15213, USA

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### ABSTRACT

Low energy ion scattering spectroscopy (LEISS) has been used to characterize the evolution of ordered structures of S on the Pd(1 1 1) surface during annealing. During exposure of the Pd(1 1 1) surface to 0.7 L H<sub>2</sub>S at 300 K—conditions that produce the  $S(\sqrt{3} \times \sqrt{3})R30$  overlayer—the intensity of the Pd LEIS signal decreases and a feature assigned to adsorbed S appears as the adsorbed layer forms. When the surface is held at 300 K after exposure to H<sub>2</sub>S is stopped, the LEIS Pd intensity partially recovers and the S signal weakens, presumably as surface S atoms assume their equilibrium positions in the  $S(\sqrt{3} \times \sqrt{3})R30$  overlayer. Subsequent annealing of the  $S(\sqrt{3} \times \sqrt{3})R30$  structure at 700 K causes it to convert into a  $S(\sqrt{7} \times \sqrt{7})R19$  overlayer, whose LEIS spectrum is identical to that of clean Pd(1 1 1). The absence of LEIS evidence for S atoms at the exposed surface of the  $S(\sqrt{7} \times \sqrt{7})R19$  overlayer is at odds with published models of a mixed Pd–S top layer. Despite the similarity of the LEIS spectra of Pd(1 1 1) and Pd(1 1 1)– $S(\sqrt{7} \times \sqrt{7})R19$ , their activities for dissociative hydrogen adsorption are very different—the former readily adsorbs hydrogen at 100 K, while the latter does not—suggesting that S exerts its influence on surface chemistry from subsurface locations.

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### 1. Introduction

Palladium is used as a catalyst for a number of important hydrogenation reactions in fine chemicals processing and petroleum refining [1–3] and as a component of dense alloy membranes used for hydrogen purification [4–7]. Exposure to sulfur-containing compounds is well known to deactivate Pd catalysts [8,9] and purification membranes [10,11], and thus, the interaction of sulfur with Pd surfaces has received significant attention from both experimental and computational researchers.

On a Pd(1 1 1) surface, two ordered overlayers of atomic sulfur are commonly formed. A ( $\sqrt{3} \times \sqrt{3}$ )R30 structure is observed when the Pd(1 1 1) surface is exposed to either H<sub>2</sub>S [12–16] or S<sub>2</sub> [17] at approximately room temperature. Experimental [12–14] and computational [18] studies conclude that S atoms in the ( $\sqrt{3} \times \sqrt{3}$ )R30 structure occupy threefold hollow fcc sites at the surface, with no S atoms in sub-surface layers. The local coverage of S in this layer is 1/3 monolayers. Upon annealing at temperatures in the range 400–700 K, a new ( $\sqrt{7} \times \sqrt{7}$ )R19 overlayer structure appears [12–14,16]. There have been several proposals for the surface structure of the ( $\sqrt{7} \times \sqrt{7}$ )R19 overlayer, based on both experiment [12,13,15,16,19] and computational modeling [20]. While there is significant diversity among the details of proposed structures, most reports agree that S does not exist in a simple adsorbed layer. Instead, S appears to penetrate the metal lattice, causing it to reconstruct to form a mixed Pd–S top-layer, which has sometimes been characterized as a two-dimensional surface sulfide [13]. The models of this structure suggest a local S coverage of either 3/7 [12,13,20] or 2/7 [15] monolayers.

In this work we have used low energy  $He^+$  ion scattering spectroscopy (LEISS) to characterize the two ordered S structures on the Pd(1 1 1) surface and to monitor their evolution during heating.

#### 2. Experimental procedures

Experiments were performed in a stainless steel ultra-high vacuum chamber with a base pressure of  $1 \times 10^{-10}$  Torr. A Pd(1 1 1) single crystal, ~10 mm diameter × 1 mm thick, was attached to an *x*, *y*, *z*,  $\theta$  manipulator by 0.63 mm Ta lead wires spot welded to its edges. The Ta wires provide both mechanical support and electrical/thermal contact to the sample, allowing control of its temperature over the range 80–1100 K. The chamber is equipped with an X-ray source (Thermo VG Scientific), a He<sup>+</sup> ion gun (Specs IQE 12/38), an energy analyzer (Specs PHOIBOS 150MCD), a mass spectrometer detector for temperature programmed desorption (TPD) experiments, low energy electron diffraction (LEED) optics (Perkin–Elmer 15-120) and an ion gun for sputter cleaning of the sample.

We prepared the surface of the Pd(1 1 1) crystal using cycles of Ar<sup>+</sup> sputtering ( $\sim$ 15  $\mu$ A/cm<sup>2</sup> at 300 K) and annealing (800 K) until S, C and O were no longer detectable by XPS. The final treatment was always a 30 min anneal at 800 K.



<sup>\*</sup> Corresponding author. Tel.: +1 412 2689517; fax: +1 412 2687139. *E-mail address*: jbmiller@andrew.cmu.edu (J.B. Miller).

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LEED patterns were acquired at ~100 K with an electron energy of ~140 eV. The XPS experiments used Al K $\alpha$  radiation and an analyzer pass energy of 100 eV. The analyzer collected electrons photoemitted at an angle of 30° from the surface normal. The low energy ion scattering spectroscopy (LEISS) experiments used 750 eV He<sup>+</sup> ions specularly reflected from the sample at an angle of ~50° from the surface normal. Low incident ion currents of ~50 nA (~5 × 10<sup>12</sup> He<sup>+</sup>/cm<sup>2</sup> over the course of a single 20 s experiment) were used in order to minimize the damage by incident He<sup>+</sup>.

Atomic S was adsorbed onto the surface of the Pd(1 1 1) crystal, cooled to ~100 K, by leaking H<sub>2</sub>S into the chamber at  $1 \times 10^{-8}$  Torr for periods chosen to deliver exposures up to 3 L. Adsorbed H<sub>2</sub>S was subsequently decomposed by heating the sample to 300 K at a rate of 4 K/s. In a similar manner, H<sub>2</sub> temperature programmed desorption (TPD) experiments were performed by first exposing the crystal to 300 L of H<sub>2</sub> (to guarantee saturation coverage) at ~100 K. The crystal was subsequently heated to 500 K at 4 K/s while positioned within 1 mm of the mass spectrometer aperture. The desorption signal at m/q = 2 was monitored as a function of temperature during heating.

## 3. Results and discussion

Fig. 1a shows the  $(1 \times 1)$  LEED pattern recorded from the clean Pd(1 1 1) surface. A S( $\sqrt{3} \times \sqrt{3}$ )R30 overlayer (Fig. 1b) was generated by exposing the annealed sample to 0.5–1.0 L H<sub>2</sub>S at 100 K, followed by heating to 300 K. This result is consistent with other published reports [12–16]. The S/Pd XPS signal intensity ratio that corresponds to this structure is 0.009 ± 0.001. Additional incremental H<sub>2</sub>S exposures up to a total of 2.0 L did not increase the S/Pd ratio as measured by XPS.

Annealing the  $S(\sqrt{3} \times \sqrt{3})R30$  structure at 700 K induces formation of the  $S(\sqrt{7} \times \sqrt{7})R19$  overlayer (Fig. 1c). This result is also consistent with published reports [12–16]. The change in the overlayer structure is not accompanied by a detectable change in the measured S/Pd XPS signal ratio. Incremental exposures of the Pd(1 1 1)– $S(\sqrt{7} \times \sqrt{7})R19$  surface to H<sub>2</sub>S (either at 700 K, or at lower temperatures followed by annealing at 700 K) up to a total exposure of 3 L changed neither the surface composition (as measured by XPS) nor the LEED pattern.

To better understand the evolution of the ordered structures, the clean, annealed Pd(111) surface was exposed to H<sub>2</sub>S at 300 K while measuring the top-layer composition using LEIS. Spectra were collected continuously during H<sub>2</sub>S exposure; acquisition time for an individual spectrum was ~20 sec. Fig. 2 displays the spectra as a function of (cumulative) H<sub>2</sub>S exposure. Each spectrum is dominated by a single feature located at approximately 705 eV, assigned to He<sup>+</sup> scattering from surface Pd atoms. The inset shows an enlargement of the S region of the same spectra, scaled to the



**Fig. 2.** LEIS spectra of the Pd(1 1 1) surface as a function of H<sub>2</sub>S exposure at 300 K. The Pd feature is located at 705 eV; the inset shows an enlargement of S feature at ~605 eV. H<sub>2</sub>S exposures are, from top to bottom, 0.0 L (clean), 0.1 L, 0.3 L, 0.5 L, and 0.7 L. Pd intensities relative to clean Pd(1 1 1) are shown at the right. As the H<sub>2</sub>S exposure increases, the Pd intensity declines and the S intensity increases.

maximum Pd signal for that spectrum. At 0.3 L and higher exposures, a small feature assignable to scattering from S atoms becomes visible at  $\sim$ 605 eV. The small size of the S signal likely reflects a low scattering cross-section-He<sup>+</sup> may prefer to neutralize, rather than to scatter, during its collision with S atoms [21]. The appearance of a S signal is accompanied by reduction of the Pd intensity, presumably as surface S atoms screen Pd atoms from incident He<sup>+</sup>. Note that Pd signal attenuation occurs at lower H<sub>2</sub>S exposures than that of the first appearance of the S feature, suggesting that Pd scattering intensity is the more sensitive measure of the presence of S atoms at the top surface. At 0.7 L H<sub>2</sub>S exposure, key features of the LEIS spectrum-appearance of the S feature at 605 eV and a reduction of Pd intensity to about 20% of the clean surface value - are the same as those observed for surfaces prepared by exposure to H<sub>2</sub>S at 100 K followed by heating to 300 K, conditions that yield the  $S(\sqrt{3} \times \sqrt{3})R30$  overlayer, with a S/Pd XPS signal ratio of  $\sim$ 0.01.

At 0.7 L,  $H_2S$  exposure was terminated, but the LEIS analysis (at 300 K) was continued for another two minutes to monitor the evolution of the overlayer. The middle three spectra displayed in Fig. 3 were acquired while holding the sample temperature at 300 K; the spectrum obtained immediately following the 0.7 L  $H_2S$  exposure from Fig. 2 is included for reference. As time elapses, the S LEIS signal at 605 eV decreases and eventually disappears. Concurrently with the loss of the S LEIS signal, the Pd signal intensity recovers,



**Fig. 1.** LEED patterns of (a) Pd(1 1 1), sputter cleaned and annealed at 800 K; (b) the  $(\sqrt{3} \times \sqrt{3})$ R30 overlayer formed on Pd(1 1 1) exposed to ~1.0 L H<sub>2</sub>S at 100 K, followed heating to 300 K, (c) the  $(\sqrt{7} \times \sqrt{7})$ R19 overlayer formed on the sample from (b) annealed at 700 K for 5 min.



**Fig. 3.** LEIS spectra of the Pd(1 1 1)–S surface as a function of time at *after* exposure to  $H_2S$  at 300 K. The Pd feature is located at 705 eV; the inset shows an enlargement of the S feature at ~605 eV. Starting at the bottom, spectra were acquired at times of 0 (the same spectrum as the bottom of Fig. 2), 20, 60, and 110 s after the  $H_2S$  exposure was discontinued. Pd intensities relative to that of clean Pd(1 1 1) are shown at the right. As time increases, the S feature disappears and the Pd feature recovers. The top spectrum was obtained after annealing at 700 K (but returned to 300 K for LEIS analysis).

but only to ~75% of its pre-exposure value. Incomplete recovery of the Pd signal is clear evidence for the presence of top-surface S atoms, and is consistent with experimental and computational characterization of S( $\sqrt{3} \times \sqrt{3}$ )R30 overlayers. Changes in the LEIS spectra that take place during annealing at 300 K may reflect changes in the surface as adsorbed S atoms assume their equilibrium positions.

After holding the sample temperature at 300 K for 2 min, the sample was annealed at 700 K for 5 min and then returned to 300 K for LEIS analysis. The LEIS spectrum at 300 K (top of Fig. 3) is nearly identical to the LEIS spectrum from the clean Pd(1 1 1) surface (top of Fig. 2): it displays comparable Pd signal at 705 eV and no S signal at 605 eV. Immediately after the LEIS experiments, LEED revealed the presence of the  $S(\sqrt{7} \times \sqrt{7})R19$  overlayer, and the S/Pd XPS ratio was 0.010—the same as that observed after 0.5–1.0 L H<sub>2</sub>S exposures at 100 K followed by stepwise heating to 300 K (to first generate the  $S(\sqrt{3} \times \sqrt{3})R30$  surface structure) and then to 700 K.

We interpret the further increase of Pd LEIS intensity after annealing at 700 K as S leaving an adsorbed layer to penetrate into the sub-surface layer of  $Pd(1 \ 1 \ 1)$ . Migration of adsorbed sulfur into the Pd(111) sub-surface is consistent with models of the  $S(\sqrt{7} \times \sqrt{7})R19$  surface that have been proposed by several researchers, which suggest a mixed top-surface that consists of both S and Pd atoms [12,13,15,16,19,20]. However, our observation that the LEIS spectra of the S(  $\sqrt{7} \times \sqrt{7}$  )R19 overlayer and the clean Pd(111) surface are nearly identical is an important-and unexpected-result. Given constant experimental parameters, LEIS signals depend primarily upon the atomic densities of surface species and their scattering cross-sections [22]. Assuming that there is no change in Pd and S scattering cross-sections as a result of adsorption on the surface, our results suggest that the top surface of the S( $\sqrt{7} \times \sqrt{7}$ )R19 overlayer contains no S atoms–or, at most, S atoms that are well-shielded by neighboring Pd atomsand that the surface density of Pd atoms is comparable to that of  $Pd(1 \ 1 \ 1)$ . In other words, we may be observing deeper penetration of S into the Pd matrix than accounted for in the mixed-top layer models of the S( $\sqrt{7} \times \sqrt{7}$ )R19 overlayer structure.



**Fig. 4.** H<sub>2</sub> TPD from the clean Pd(111) surface (top) and the Pd(111)– $S(\sqrt{7} \times \sqrt{7})R19$  surface (bottom), obtained by heating at 4 K/s while monitoring the mass spectrometer signal at m/q = 2. Exposure of 300 L H<sub>2</sub> at 100 K was sufficient to saturate the clean Pd(111) surface, but, under the same conditions of exposure, the Pd(111)– $S(\sqrt{7} \times \sqrt{7})R19$  surface did not adsorb significant H<sub>2</sub>.

The view that S atoms on Pd(1 1 1) have been displaced into the immediate sub-surface layer by annealing to 700 K is balanced, at least partially, by the XPS results. The disappearance of LEIS evidence of top-surface S that takes place as the ordered  $S(\sqrt{3} \times \sqrt{3})$ R30 surface converts to the  $S(\sqrt{7} \times \sqrt{7})$ R19 structure is not accompanied by a detectable decrease in the XPS S/Pd signal ratio. Thus, XPS suggests that the extent of penetration of S into the near surface is not great enough to cause significant attenuation of the S photoelectron signal by Pd.

The similarity between the LEIS spectra of the  $S(\sqrt{7} \times \sqrt{7})R19$ overlayer and the clean Pd(111) surface also suggests that the two might have similar surface chemistries. We compared TPD spectra of H<sub>2</sub> from the two surfaces; results appear in Fig. 4. As expected [23–26], the clean Pd surface readily adsorbs H<sub>2</sub> at 100 K and, during heating, H<sub>2</sub> subsequently desorbs with a peak maximum in the vicinity of 300 K. The Pd(1 1 1)–S( $\sqrt{7} \times \sqrt{7}$ )R19 surface, in contrast, does not adsorb H<sub>2</sub> under the same conditions of exposure. This result is consistent with those reported for the Pd(110) surface, on which the hydrogen adsorption capacity declines linearly with increasing S coverage, resulting in complete deactivation of the surface at a S coverage of  $\theta_{\rm S}$  = 0.28 ML [26]. Because the  $S(\sqrt{7} \times \sqrt{7})R19$  overlayer adsorbs no H<sub>2</sub>, we conclude that either S in the subsurface exerts an electronic effect on the Pd or that structural features of surface sites that are necessary for the dissociative adsorption of H<sub>2</sub> are not adequately probed by LEIS.

### 4. Conclusions

Exposure of a Pd(1 1 1) surface to H<sub>2</sub>S at conditions that produce the  $S(\sqrt{3} \times \sqrt{3})R30$  structure–0.5–1.0 L with heating to 300 K–results in decreasing Pd LEIS intensity and appearance of a new feature assigned to surface S. During continued heating at 300 K after exposure to H<sub>2</sub>S, the LEIS Pd intensity partially recovers and the S signal disappears, presumably as S atoms assume their equilibrium positions in the  $S(\sqrt{3} \times \sqrt{3})R30$  structure at 700 K causes it to convert into a  $S(\sqrt{7} \times \sqrt{7})R19$  overlayer, whose LEIS spectrum is identical to

that of clean Pd(1 1 1). The absence of LEIS evidence for top-surface S in the S( $\sqrt{7} \times \sqrt{7}$ )R19 overlayer is at odds with published models of a mixed Pd and S top surface.

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#### References

- G.J. Antos, A.M. Aitani, J.M. Parera, Catalytic Naphtha Reforming: Science and Technology, Marcel Dekker, New York, 1995.
- [2] H. Arnold, F. Dobert, J. Gaube, in: G. Ertl, H. Knozinger, J. Weitkamp (Eds.), Handbook of Heterogeneous Catalysis, (Wiley–VCH, Weinheim, Germany, 1997, p. 1211.
- [3] A. Sarkany, Appl. Catal. A 165 (1997).
- [4] R.E. Buxbaum, T.L. Marker, J. Membrane Sci. 85 (1993) 29.
- [5] B.H. Howard et al., J. Membrane Sci. 241 (2004) 207.
- [6] P. Kamakoti et al., Science 307 (2005) 569.

- [7] B.D. Morreale et al., J. Membrane Sci. 212 (2003) 87.
- [8] J. Oudar, H. Wise, Deactivation and Poisoning of Catalysts, Dekker, New York, 1991.
- [9] J.A. Rodriguez, J. Hrbek, Acc. Chem. Res. 32 (1999) 719.
- [10] B.D. Morreale, University of Pittsburgh, 2006.
- [11] B.D. Morreale et al., J. Membrane Sci. 241 (2004) 219.
- [12] J.G. Forbes, A.J. Gellman, J.C. Dunphy, M. Salmeron, Surf. Sci. 279 (1992) 68.
- [13] M.E. Grillo, C. Stampfl, W. Berndt, Surf. Sci. 317 (1994) 84.
- [14] F. Maca, M. Scheffler, W. Berndt, Surf. Sci. 160 (1985) 467.
- [15] S. Speller, T. Rauch, J. Bommermann, P. Borrmann, W. Heiland, Surf. Sci. 441 (1999) 107.
- [16] V.R. Dhanak, A.G. Shard, B.C.C. Cowie, A. Santoni, Surf. Sci. 410 (1998) 321.
- [17] C.H. Patterson, R.M. Lambert, Surf. Sci. 187 (1987) 339.
- [18] D.R. Alfonso, Surf. Sci. 596 (2005) 229.
- [19] W. Liu, K.A.R. Mitchell, W. Berndt, Surf. Sci. Lett. 393 (1997) L119.
- [20] D.R. Alfonso, Surf. Sci. 601 (2007) 4899.
- [21] Y. Luo, M. Han, D.A. Slater, R.M. Osgood, J. Vac. Sci. Tech. A 18 (2000) 438.
- [22] H.H. Brongersma, M. Draxler, M. de Ridder, P. Bauer, Surf. Sci. Rep. 62 (2007) 63.
- [23] M.G. Cattania, V. Penka, R.J. Behm, K. Christmann, G. Ertl, Surf. Sci. 126 (1983) 382.
- [24] U. Muschiol, P.K. Schmidt, K. Christmann, Surf. Sci. 395 (1998) 182.
- [25] H. Okuyama, W. Siga, N. Takagi, M. Nishijima, T. Aruga, Surf. Sci. 401 (1998) 344.
- [26] M.L. Burke, R.J. Madix, Surf. Sci. 237 (1990) 1.